

Taking an average is a linear process.

- For a discrete variable j and probability P_j

$$\langle \alpha f(j) + \beta g(j) \rangle = \alpha \langle f(j) \rangle + \beta \langle g(j) \rangle$$

- Similarly for a continuous variable x and probability density ftn $\rho(x)$

$$\langle \alpha f(x) + \beta g(x) \rangle = \alpha \langle f(x) \rangle + \beta \langle g(x) \rangle$$

proof.)

$\alpha, \beta = \text{constants}$
(i.e. indep of j or x)

Only the discrete variable case is discussed here, as the proof for the other case is very similar ("left for homework").

$$\begin{aligned} \langle \alpha f(j) + \beta g(j) \rangle &= \sum_j \{ \alpha f(j) + \beta g(j) \} P_j \\ &= \alpha \sum_j f(j) P_j + \beta \sum_j g(j) P_j \\ &= \alpha \langle f(j) \rangle + \beta \langle g(j) \rangle \end{aligned}$$

Math Trick

M①

① What is $\int_{-\infty}^{\infty} dx e^{-x^2}$? $\sqrt{\pi}$

② What is $\int_{-\infty}^{\infty} dx e^{-\lambda x^2}$? $\sqrt{\frac{\pi}{\lambda}}$

③ What is $\int_{-\infty}^{\infty} dx x^{2n} e^{-x^2}$? $\sqrt{\pi} \cdot \frac{1}{2} \cdot \frac{3}{2} \cdots \frac{2n-1}{2}$
($n \geq 1$)

① $\left(\int_{-\infty}^{\infty} dx e^{-x^2} \right)^2 = \int_{-\infty}^{\infty} dx \int_{-\infty}^{\infty} dy e^{-x^2} e^{-y^2} = \int_{-\infty}^{\infty} dx \int_{-\infty}^{\infty} dy e^{-(x^2+y^2)} = *$

$\begin{cases} x = r \cos \theta \\ y = r \sin \theta \end{cases} \quad (r^2 = t)$
 $* = \int_0^{\infty} dr \int_0^{2\pi} d\theta r e^{-r^2} = 2\pi \int_0^{\infty} dr r e^{-r^2} = \pi \int_0^{\infty} dt e^{-t} = \pi$
QED

② $\int_{-\infty}^{\infty} dx e^{-\lambda x^2} = \frac{1}{\sqrt{\lambda}} \int_{-\infty}^{\infty} dx' e^{-x'^2} = \sqrt{\frac{\pi}{\lambda}}$ by ①

③ $\int_{-\infty}^{\infty} dx x^{2n} e^{-x^2} = \left(-\frac{\partial}{\partial \lambda} \right)^n \int_{-\infty}^{\infty} dx e^{-\lambda x^2} \Big|_{\lambda=1} = \left(-\frac{\partial}{\partial \lambda} \right)^n \left(\sqrt{\frac{\pi}{\lambda}} \cdot \lambda^{-\frac{1}{2}} \right) \Big|_{\lambda=1}$
 $= \sqrt{\pi} \cdot \underbrace{\frac{1}{2} \cdot \frac{3}{2} \cdot \frac{5}{2} \cdots \frac{2n-1}{2}}_n \left[= \sqrt{\pi} \cdot \frac{(2n)!}{n! 2^{2n}} \right]$

Math Info

~~Q1~~ Gaussian distribution ftn

$$f(x) = A e^{-\frac{x^2}{2\sigma^2}}$$

Normalize: $A \int_{-\infty}^{\infty} dx e^{-\frac{x^2}{2\sigma^2}} = 1$

$$A \sqrt{2\sigma} \sqrt{\pi} = 1 \quad \therefore A = \frac{1}{\sigma\sqrt{2\pi}}$$

Mean value: $\langle x \rangle = 0$

~~Q2~~

$$\begin{aligned} \langle x^2 \rangle &= \int_{-\infty}^{\infty} x^2 f(x) dx = A \int_{-\infty}^{\infty} x^2 e^{-\frac{x^2}{2\sigma^2}} dx \\ &= A \cdot (\sqrt{2\sigma})^3 \cdot \frac{1}{2} \cdot \sqrt{\pi} = A \cdot \sqrt{2} \sigma^3 \sqrt{\pi} = \sigma^2 \end{aligned}$$

$$\therefore \text{Standard dev.} = \sqrt{\langle x^2 \rangle - \langle x \rangle^2} = \sigma$$

Gaussian Distribution Function

$$f(x) = \frac{1}{\sqrt{2\pi} \cdot \sigma} e^{-\frac{x^2}{2\sigma^2}}$$

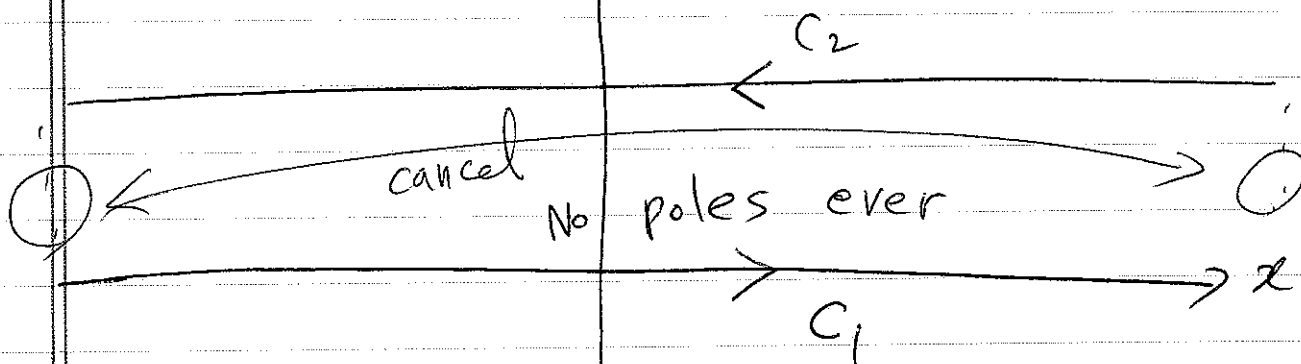
However, warning: do not confuse wave function $\Psi(x)$ and $\rho(x)$.
For $\rho(x)$ to be the normalized gaussian ftn given above
 $\Psi(x) = \sqrt{f(x)}$ up to a phase factor.

Math Trick

$$\int_{-\infty}^{\infty} e^{-(ax^2+bx)} dx = \int_{-\infty}^{\infty} dx e^{-a\left(x+\frac{b}{2a}\right)^2} e^{\frac{b^2}{4a}}$$

$$= \sqrt{\frac{\pi}{a}} \cdot e^{\frac{b^2}{4a}}$$

Even if $a, b \in \mathbb{Z}$, this holds.
 (Even if $a, b \in \mathbb{C}$, this holds.)
 $(\operatorname{Re}\{a(x+\frac{b}{2a})^2\} > 0)$



$$\oint_{C_1} e^{-az^2} = \oint_{C_2} e^{-az^2}$$

from Cauchy theorem
of complex analysis.

You are not required to know this.
 If you know the theorem, good for you.
 If you don't, please take my word for it.
 Ref. Arfken etc.

Math note

$$\psi_n = \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi x}{a}\right), \quad n = \text{integer} \rightarrow 0$$

$$\int_0^a dx \psi_m^* \psi_n = \delta_{mn} = \begin{cases} 1 & m=n \\ 0 & m \neq n \end{cases}$$

proof?

$$I_{mn} = \int_0^a dx \psi_m^* \psi_n = \frac{2}{a} \int_0^a dx \sin\left(\frac{m\pi x}{a}\right) \sin\left(\frac{n\pi x}{a}\right)$$

$$= \frac{a}{\pi} \cdot \frac{2}{a} \int_0^\pi dy \sin(my) \sin(ny)$$

$$\uparrow$$

$$\frac{\pi x}{a} = y$$

Now use $\sin \theta = \frac{e^{i\theta} - e^{-i\theta}}{2i}$

$$I_{mn} = \frac{2}{\pi} \int_0^\pi dy \frac{(e^{imy} - e^{-imy})(e^{iny} - e^{-iny})}{-4}$$

(\leftarrow Euler relation $e^{i\theta} = \cos\theta + i\sin\theta$)

$$= -\frac{1}{2\pi} \int_0^\pi dy \left\{ e^{i(m+n)y} + e^{-i(m+n)y} - e^{i(m-n)y} - e^{-i(m-n)y} \right\}$$

$$= -\frac{1}{2\pi} \left(\frac{((-1)^{m+n} - 1) - ((-1)^{m+n} - 1)}{i(m+n)} - \frac{((-1)^{m-n} - 1) - ((-1)^{m-n} - 1)}{i(m-n)} \right)$$

same thing

$$= 0 \quad \text{as long as } m-n \neq 0$$

$$m+n \neq 0 \leftarrow \text{impossible}$$

If $m=n$, $I_{mn} = 1$

QED.

Math note. What in the world does it mean?

$$\textcircled{1} \quad \frac{1}{2\pi} \int_{-\infty}^{\infty} dk e^{ikx} = \delta(x)$$

$$\textcircled{2} \quad \text{or, } \frac{1}{2\pi} \int_{-\infty}^{\infty} dx e^{ikx} = \delta(k)$$

① is completely equivalent to ②.

Will discuss ②.

② is quite troublesome indeed. When $k \neq 0$, it is clear that $\frac{1}{2\pi} \int_{-\infty}^{\infty} dx e^{ikx}$ has many periods

of waves that are compensating each other, but why is it necessarily 0? Depending on how one takes limits $\pm\infty$, it can be a finite number!

When $k=0$, clearly $\frac{1}{2\pi} \int_{-\infty}^{\infty} dx e^{ikx} \rightarrow \infty!$

A mathematical way to better define the above expression is $\frac{1}{2\pi} \int_{-\infty}^{\infty} dx e^{i(k \pm i\epsilon)x} \equiv f(k)$

$$f(k) = \frac{1}{2\pi} \left[\int_{-\infty}^0 dx e^{i(k-i\epsilon)x} + \int_0^{\infty} dx e^{i(k+i\epsilon)x} \right]$$

$$= \frac{1}{2\pi} \left[\frac{1 - e^{i(k-i\epsilon)\infty}}{i(k-i\epsilon)} + \frac{e^{i(k+i\epsilon)\infty} - 1}{i(k+i\epsilon)} \right] = \frac{1}{2\pi i} \left(\frac{1}{k-i\epsilon} - \frac{1}{k+i\epsilon} \right)$$

See comment in M③.

← From Cauchy's theorem of complex analysis, can show

$$\int_{-\infty}^{\infty} f(k) dk = 1$$

With $\epsilon > 0$, but ^{much} smaller than any finite k , $\epsilon \ll |k|$

$$f(k) = 0 \text{ for any finite } k \rightarrow \boxed{\therefore f(k) = \delta(k)}$$

Dirac delta function $\delta(x)$

Definition: $\delta(x) = 0, x \neq 0$ and $\int_{-\infty}^{\infty} dx \delta(x) = 1$

There are various ways to ~~derive~~ derive $\delta(x)$
 [Can show all of them satisfy the above def.]

① $\delta(x) = \theta'(x)$ where $\theta(x) = \begin{cases} 1 & x \geq 0 \\ 0 & x \leq 0 \end{cases}$ \Rightarrow equal sign can be used at either of the two or $\theta(0) = \frac{1}{2}$
 $\Rightarrow \delta(x) = (\theta(x) + \text{const.})'$

② $\delta(x) = \lim_{\sigma \rightarrow 0} \frac{1}{\sqrt{2\pi\sigma^2}} \exp(-\frac{x^2}{2\sigma^2})$ \leftarrow normalized gaussian

③ $\delta(x) = \lim_{a \rightarrow 0} \frac{a}{\pi} \frac{1}{x^2 + a^2}$ \leftarrow normalized lorentzian

④ $\delta(x) = \lim_{a \rightarrow 0} \begin{cases} \frac{1}{2a} & -a \leq x \leq a \\ 0 & \text{elsewhere} \end{cases}$ \leftarrow normalized square pulse

⑤ $\delta(x) = \lim_{n \rightarrow \infty} \frac{\sin nx}{\pi x}$ \leftarrow These two are related since $\frac{1}{2\pi} \int_{-n}^n e^{ikx} dk = \frac{\sin nx}{\pi x}$

⑥ $\delta(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} dk e^{ikx}$

②-⑤ can be generalized as $\delta(x) = \text{limit of any normalized ftn } \int F(x) dx = 1$ with "width, W " of $F(x)$ going to 0.

①, ⑥ were explained before. $\delta(x) = \lim_{W \rightarrow 0} F(x)$

Properties of $\delta(x)$ proof "left for readers"

- ① ~~$f(x) \delta(x) = f(0) \delta(x)$~~ $f(x) \delta(x) = f(0) \delta(x)$, ② $\delta(-x) = \delta(x)$
- ③ $\delta(ax+b) = \frac{1}{|a|} \delta(x + \frac{b}{a})$, ④ $\delta(f(x)) = \left| f'(x) \right|^{-1} \delta(x - x_0)$ $f(x_0) = 0$

Note $\int_0^{\infty} \frac{\sin x}{x} dx = \frac{\pi}{2}$